

# Quantum statistics and interactions

Exercise session II - C. Winkelmann - UGA/Phelma

## Introduction to Quantum Fields

Let  $|0\rangle$  be the vacuum state,  $\{|\phi_i\rangle\}$  an orthonormal basis of single particle states, and  $|\mathbf{r}\rangle$  the delta-shape wave function of a particle localized at position  $\mathbf{r}$ .

### 1 Single-particle wave function

#### 1.1

Prove the following orthonormality and closure relations

$$\int d^3\mathbf{r} \phi_i^*(\mathbf{r})\phi_j(\mathbf{r}) = \delta_{ij}, \quad \sum_i \phi_i^*(\mathbf{r})\phi_i(\mathbf{r}') = \delta(\mathbf{r} - \mathbf{r}').$$

#### 1.2

We define  $\Psi^\dagger(\mathbf{r})$  as the operator creating a particle at position  $\mathbf{r}$ . Thus we can write  $\Psi^\dagger(\mathbf{r})|0\rangle = |\mathbf{r}\rangle$ . Show that  $\Psi^\dagger(\mathbf{r})|0\rangle = \sum_i \phi_i^*(\mathbf{r}) a_i^\dagger|0\rangle$ .

#### 1.3

We can thus define the field operator associated to this basis,  $\Psi^\dagger(\mathbf{r}) = \sum_i \phi_i^*(\mathbf{r}) a_i^\dagger$ . Give also the adjoint operator  $\Psi(\mathbf{r})$ .

#### 1.4

We consider the case of a spinless boson. Show that  $[\Psi(\mathbf{r}), \Psi^\dagger(\mathbf{r}')] = \delta(\mathbf{r} - \mathbf{r}')$  and  $[\Psi(\mathbf{r}), \Psi(\mathbf{r}')] = 0$ .

#### 1.5

Show that

$$N = \int d^3\mathbf{r} \Psi^\dagger(\mathbf{r})\Psi(\mathbf{r}),$$

and give an interpretation of  $\Psi^\dagger(\mathbf{r})\Psi(\mathbf{r})$ .

#### 1.6

We now consider the case of a spin-1/2 fermion and include the spin degree of freedom as an additional quantum number  $\sigma$ , writing now the operators  $a_{i,\sigma}$  and  $\Psi_\sigma(\mathbf{r})$ . Write the relevant fermionic anticommutation relations, such as between  $a_{i,\sigma}$  and  $a_{j,\sigma'}^\dagger$ .

#### 1.7

We now define the field operators as

$$\Psi_\sigma^\dagger(\mathbf{r})|0\rangle = \sum_i \phi_i^*(\mathbf{r}) a_{i,\sigma}^\dagger|0\rangle.$$

Show that  $\{\Psi_\sigma(\mathbf{r}), \Psi_{\sigma'}^\dagger(\mathbf{r}')\} = \delta(\mathbf{r} - \mathbf{r}') \delta_{\sigma,\sigma'}$  and  $\{\Psi_\sigma(\mathbf{r}), \Psi_{\sigma'}(\mathbf{r}')\} = 0$ .

## 1.8

Define the number operator  $N$  in this case.

## 1.9

We consider the specific example of electrons in a cubic box of volume  $V = L^3$ . We use plane waves as a basis of wavefunctions

$$\phi_{\mathbf{k}}(\mathbf{r}) = \langle \mathbf{r} | \mathbf{k} \rangle = \frac{e^{i\mathbf{k} \cdot \mathbf{r}}}{L^{3/2}}.$$

Give the expressions of  $\Psi_{\sigma}(\mathbf{r})$  and  $\Psi_{\sigma}^{\dagger}(\mathbf{r})$ , and interpret. Eventually show that the total number operator is given by

$$N = \sum_{\mathbf{k}, \sigma} a_{\mathbf{k}, \sigma}^{\dagger} a_{\mathbf{k}, \sigma}.$$

## 2 Many-particle wave function

As above, and referring to the same single-particle basis, we can construct the many-particle wave function from the vacuum state by defining

$$|\mathbf{r}_1, \mathbf{r}_2, \dots\rangle := \Psi^{\dagger}(\mathbf{r}_1) \Psi^{\dagger}(\mathbf{r}_2) \dots \Psi^{\dagger}(\mathbf{r}_N) |0\rangle.$$

In case the particles have an internal degree of freedom  $\sigma$ , like the spin, this can be written

$$|\mathbf{r}_1 \sigma_1, \mathbf{r}_2 \sigma_2, \dots\rangle := \Psi_{\sigma_1}^{\dagger}(\mathbf{r}_1) \Psi_{\sigma_2}^{\dagger}(\mathbf{r}_2) \dots \Psi_{\sigma_N}^{\dagger}(\mathbf{r}_N) |0\rangle. \quad (1)$$

### 2.1

Write these bras in ket form.

### 2.2

Still considering the same single-particle basis as discussed previously, we construct the many-particle state

$$a_{1, \tau_1}^{\dagger} a_{2, \tau_2}^{\dagger} \dots a_{N, \tau_N}^{\dagger} |0\rangle := |1 \tau_1, 2 \tau_2, \dots, N \tau_N\rangle. \quad (2)$$

Show that  $\langle \mathbf{r}_i \sigma_i | j \tau_j \rangle = \phi_j(\mathbf{r}_i) \delta_{\sigma_i, \tau_j}$ .

### 2.3

Write the scalar product of the two vectors defined in (1) and (2), respectively.

### 2.4

Wick's theorem (not proven) allows stating that if there exists a permutation  $p$  transforming  $(1, 2, \dots, N)$  into  $(i_1, i_2, \dots, i_N)$ , then

$$\langle 0 | a_{i_N, \sigma_N} \dots a_{i_2, \sigma_2} a_{i_1, \sigma_1} a_{1, \tau_1}^{\dagger} a_{2, \tau_2}^{\dagger} \dots a_{N, \tau_N}^{\dagger} |0\rangle = \text{sign}(p) \times \delta_{\sigma_1, \tau_{p(1)}} \dots \delta_{\sigma_N, \tau_{p(N)}}.$$

What would the value of the above matrix element if there was no such permutation?

### 2.5

Write the scalar product from question 2.2 in the form of a Slater determinant.