

Very low shot noise in carbon nanotubes

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Abstract. We have performed noise measurements on suspended ropes of single wall carbon nanotubes (SWNT) between 1 and 300 K for different values of dc current through the ropes. We find that the shot noise is suppressed by more than a factor 100 compared to the full shot noise $2eI$. We have also measured an individual SWNT and found a level of noise which is smaller than the minimum expected. Another finding is the very low level of $1/f$ noise, which is significantly lower than previous observations. We propose two possible interpretations for this strong shot noise reduction: i) Transport within a rope takes place through a few nearly ballistic tubes within a rope and possibly involves non integer effective charges with $e^* \sim 0.3e$. ii) A substantial fraction of the tubes conduct with a strong reduction of effective charge (by more than a factor 50).

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1 Introduction

Metallic single wall carbon nanotubes (SWNT) are known to be model systems for the study of one dimensional (1D) electronic transport. Depending on diameter and helicity, SWNT have at most 2 conducting channels and in the absence of disorder their minimum resistance is predicted to be $h/4e^2 \approx 6.5 k\Omega$ [1, 2]. In 1D, electronic interactions are expected to lead to a breakdown of the Fermi liquid state. Nanotubes should then be described by Luttinger Liquids (LL) theories [3, 4], with collective low energy excitations. Indications of the validity of LL description with repulsive interactions in SWNT were given by the measurement of a resistance diverging as a power law with temperature down to 10 K [5]. It has been shown that another experimental signature of a Luttinger liquid is the existence of non integer charge excitations $e^* = ge$ where g is the Luttinger parameter which depends on the interaction strength ($g=1$ in the absence of interactions.) These excitations can in principle be detected in a 2-wire measurement on a sample containing an impurity in the weak backscattering regime [6, 7]. Indeed, non integer charge excitations create a current noise which has the form:

$$S_I = 2e^*i_b \quad (1)$$

where i_b is the back-scattered current. Edge states in 2D electron systems in the fractional quantum Hall regime

have provided an elegant demonstration of this prediction [8]. However carbon nanotubes, contrary to the edge states in the fractional quantum hall regime, are expected to constitute a non chiral Luttinger Liquid where forward and backward electrons are not spatially separated. Theoretical predictions for the non chiral LL are quite controversial. Although the initial prediction eq.1 [7] was obtained for a non chiral infinite LL, it has been recently claimed [9, 10] that the shot noise in a LL between two reservoirs does not reveal the quasi-particle non integer charge and the authors in [10] predict instead an integer charge of $1e$, in the expression of shot noise. These calculations however rely on drastic hypothesis concerning screening of electron interactions and boundary conditions with the reservoirs which validity can be questioned [11]. The investigation of shot noise in carbon nanotubes is therefore an important issue. The measurements we present on SWNT show a surprisingly high shot noise reduction.

2 Experimental set-up

The SWNT are prepared by an electrical arc method with a mixture of nickel and yttrium as a catalyst [12, 13]. SWNT with diameters of the order of 1.4 nm are obtained. They are purified by the cross-flow filtration method [13]. The tubes usually come assembled in ropes of a few hun-

dred parallel tubes, but individual tubes can also be obtained after chemical treatment with a surfactant [14]. Isolation of an individual rope and connection to measuring pads are performed according to the following nano-soldering technique. A target covered with nanotubes is placed above a metal-coated [16] suspended Si_3N_4 membrane, in which a roughly $1 \mu m$ by $100 \mu m$ slit has been etched with a focused ion beam. Following ref. [15] a focused UV laser beam pulse (power 10 kW) is applied for 10 ns, a nanotube drops from the target and connects the edges of the slit. Since the metal electrodes on each side of the slit are locally molten, the tube gets soldered into good contact. The resulting samples can then be characterized using transmission electronic microscopy. The measurement of large supercurrents when the tubes are soldered to superconducting electrodes [17] indicate that the contact resistance is much smaller than the intrinsic resistance of the tubes.

Using this technique, we have obtained individual carbon nanotubes such as the sample presented in fig.(5), with resistances as low as $10 k\Omega$, that present a quasi metallic behavior down to very low temperature with less than a 25% increase of resistance between 300 K and 1 K. Ropes containing few tens of nearly ordered carbon nanotubes soldered using the same technique present a wide range of measured resistances that can vary between 100 and $10^5 \Omega$ at 300 K. These ropes are generally metallic when their resistances are below $10 k\Omega$ at room temperature. The data presented here are taken on ropes that have a resistance less than $1 k\Omega$. (More resistive ropes usually had a level of $1/f$ noise that prevents the measurement of shot noise.) The temperature dependence of these low resistance ropes (see inset of fig. 1) is also very weak but in contrast with what is observed in individual tubes, it is not monotonous as already reported in [18]. The resistance decreases linearly as temperature is lowered between room temperature and 30 K indicating the freezing-out of phonon modes, and then increases as T is further decreased, just as in individual tubes.

The noise was recorded as a voltage fluctuation $\Delta V(t)$ across the samples biased with various DC currents I_{DC} . The spectrum of voltage noise power

$P_V(f) = 2 \int \overline{\Delta V(t)\Delta V(O)} \exp(2i\pi ft) dt$ was averaged during over roughly 10 minutes and converted into a current noise according to $P_I = h.P_V/R^2$, where R is the sample resistance and $h = (1 + R/Z)^2$ is a correction factor which accounts for the finite impedance $Z \gg R$ of the biasing circuit. For all the samples tested, this correction factor never exceeds the value 1.2. In practice, a spectrum analyser calculated the noise power spectrum by correlating the output signals from two low noise differential preamplifiers (model LI-75A by NF Electronic Instruments) which were independently connected across the samples. The measuring procedure has been checked by replacing the samples with a 560Ω macroscopic resistor, for which zero shot noise is expected. The experiment, conducted at 4.2 K, gave the expected result within $\pm 2eI/2000$ (see fig. 2). Another validation of the apparatus and correlation procedure was provided by the mea-

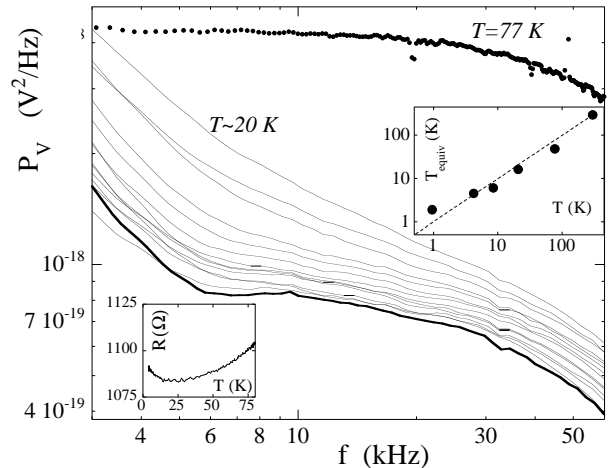


Fig. 1. Typical voltage noise power spectra P_V at 77 K (zero current bias) and 20.6 K ($I_{DC} = 0$ (thick line), 20, 40, ..., 260 μA) for the Rope A. The high frequency cut-off is due to a protective low pass filtering. The dashed lines are the expected Johnson Nyquist noise floor. Lower Inset: Temperature dependence of the resistance of Rope A. Upper Inset: Noise floor for $I_{DC} = 0$ in K ($P_V/4kR$) versus temperature.

surement of the full shot noise $2eI$ across the PN junction of a commercial transistor at room temperature.

3 Results

Typical noise spectra obtained on a rope of 200 tubes, which resistance varies between 1085Ω and 1100Ω are shown on fig.1 for I_{DC} currents up to 260 μA (Rope A). On all spectra, we have eliminated parasitic peaks (due to harmonics of 50 Hz and radiations of monitors and screens), which are independent of I_{DC} . They represent about 1% of the investigated frequencies. To increase the readability of the fig.1, the data used for the plot have been smoothed over a 600 Hz frequency window. One can note also the high frequency cutoff above 30 kHz which is due to input filters in the experiment.

Every sample exhibits three intrinsic noise contributions: Johnson Nyquist thermal noise, shot noise and low frequency noise. In addition, the major sources of background noise are the circuit resistances and preamplifiers voltage and current noise. These different parasitic contributions have been characterized independently with specific experiments, in particular at 140 mK in order to reduce the thermal noises. For all measurements, the total background noise is dominated by the amplifiers current noise ($130 fA.Hz^{-0.5}$ at 10 kHz) and it always remains smaller than the noise from the sample itself.

For $I_{DC} = 0$ and beyond the low frequency region, the white noise P_I^0 is identified with Johnson Nyquist noise. Its level was determined after correction for the high frequency cutoff, which is due to known filter capacitances. This noise power can be expressed as an equivalent tem-

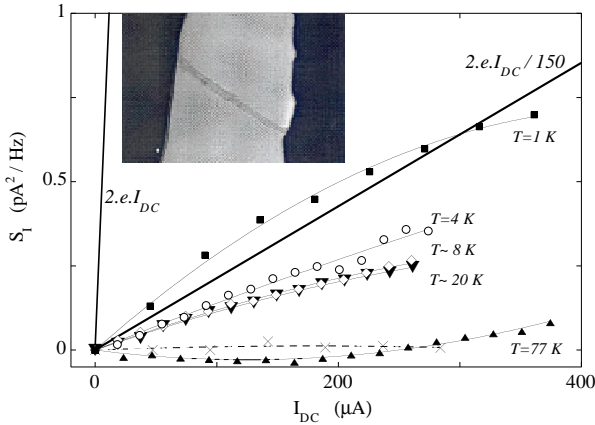


Fig. 2. Shot noise $S_I = P_I(I_{DC}) - P_I(0)$ versus bias current I_{DC} for the Rope A at different temperatures. The thick solid lines represent the 1/150 reduced and full shot noise $2eI_{DC}$. X : shot noise of a 560Ω macroscopic resistor at 4.2 K (control sample). Inset : TEM micrograph of the Rope A, from which we deduce its length $L = 0.4 \mu m$ and the number of tubes N from the measured diameter D through $N = (D/(d+e))^2$, where d is the diameter of a single tube ($d=1.4 \text{ nm}$), and e is the typical distance between tubes in a rope ($e=0.2 \text{ nm}$).

perature T_{eq} according to $T_{eq} = P_I^0 R / 4k_B$. The upper inset of fig.1 shows that T_{eq} is always close to the measured value of T , which provides a basic experimental check.

The quantity we consider from now on is the difference between the noise power spectra obtained with and without DC current bias : $P_I(I_{DC}) - P_I(0)$. This difference only contains shot noise and low frequency noise. We find that the low frequency noise roughly scales as I^2/f , which is the basic signature of $1/f$ noise originating from the a population of active two-level systems in a metallic conductor [19].

Note that the amplitude of $1/f$ noise we measure is very small compared to the values already published on carbon nanotubes [21,22]. A possible explanation is the suspended character of our samples in contrast with deposited nanotubes which can interact with residual mobile charges present on any substrate. Following the analysis in [22] we characterize the amplitude of $1/f$ noise by the dimensionless coefficient B defined as $S_I/I_{dc}^2 = B/f$. For rope A, we find $B \simeq 10^{-13}$ nearly independent of temperature between 300 K and 1 K. For the highest currents (typically $I_{DC} > 350 \mu A$), an additional low frequency noise appears with a $1/f^2$ dependence which saturates at very low frequency. This behavior is typical of telegraphic noise which characteristic frequency increases with I_{DC} (see fig. 1) and is similar to previous observations in nanometer scales metallic samples where the resistance noise is dominated by the activity of an individual or a very few number of two-level-systems.[20]. We fit the tail of the low frequency noise by a $1/f^\alpha$ power law, and the remaining white noise is identified to the shot noise S_I . Above 5-10 kHz , the amplitude of the white noise level is not found to depend on the details of the low frequency fit, at least for α arbitrarily chosen between 1 and 2.

We now focus exclusively on this shot noise. In fig. 2, S_I is plotted as a function of I_{DC} for Rope A. As expected we observe a roughly linear dependence. Interestingly, the current noise is strongly reduced (by more than a factor 150) compared to the full shot noise value $S_I = 2eI$. For this reason it cannot be detected for currents below $10 \mu A$. This strong reduction of shot noise is the central result of our paper. Note however that we cannot discriminate in these experiments the intrinsic noise of the tubes from the noise of the contacts. It is thus possible that the intrinsic noise of the nanotubes is even smaller than these measured values.

This small shot noise amplitude decreases systematically as temperature is increased between 1 and 20 K and is not measurable anymore at 77 K just like in a macroscopic resistor. Shot noise is indeed expected to disappear at temperatures for which inelastic electron phonon scattering takes place inside the sample, with scattering lengths much smaller than the sample length. (And indeed the increase of resistance measured above 30 K indicates that e-phonon scattering takes place inside the sample at this temperature).

Similar shot noise reductions were also seen (see fig.3) in the less resistive ropes B ($R = 495 \Omega$) and C ($R = 186 \Omega$) where reduction factors are respectively measured to be 270 and 150 at 1K. We have also performed noise measurements on an individual tube at 1 K (fig. 4) which resistance varies between 10.5 and 12.5 $k\Omega$ between 300 K and 1 K. The amplitude of $1/f$ noise measured in this sample is much higher than in ropes, and presents strong non-linearities with I_{dc}^2 above $0.1 \mu A$ as shown on fig.5. At current below $0.1 \mu A$ it can be described by the coefficient $B = 10^{-9}$ which is about 10^2 smaller than the value recently measured at 8 K in an individual tube 5 times more resistive by Postma *et al.* [22]. The strong difference by more than a factor 10^4 of the amplitude of $1/f$ noise between ropes containing 300 tubes and an individual tube may seem surprising. But it can be qualitatively understood considering that transport through a rope essentially probes the most conducting tubes which are also likely to be the less noisy ones. This large amount of $1/f$ noise in the individual tube sample makes the analysis of the shot noise component more difficult. The data shown in fig. 5 shows that there is barely no detectable shot noise in this sample. We can nevertheless give an upper-bound estimate of the amount of shot noise which are the data points shown.

4 Discussion

We now discuss the possible physical mechanisms which could be at the origin of the strong shot noise reduction.

i) Strong electron phonon scattering in the sample or at the contacts. Electron phonon scattering characterized by a typical scattering length l_{eph} has been shown to induce a shot noise reduction of a factor l_{eph}/L for a sample of length L much longer than l_{eph} [24]. However in order to explain the noise reduction factor observed in the present experiments an inelastic length as short as 5 nm is needed

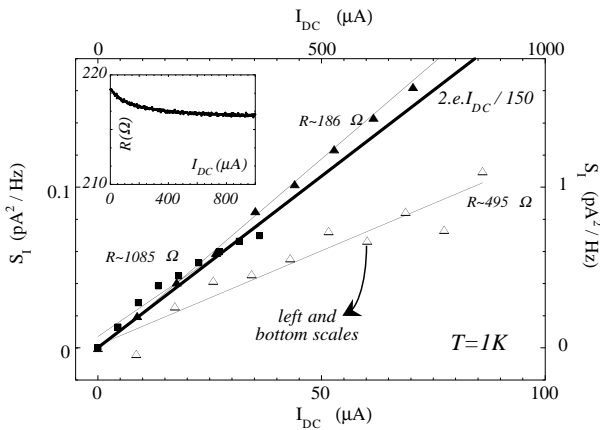


Fig. 3. Shot noise S_I versus bias current I_{DC} at 1 K for the different Ropes A (full squares), B (empty triangles) and C (full triangles) which have similar lengths and diameters. Note different scales both on Y and X axis for the rope B with $R = 495 \Omega$. The thick solid line represents the reduced shot noise $2eI_{DC}/150$ on both scales. Inset: Differential resistance versus applied dc current for the rope C.

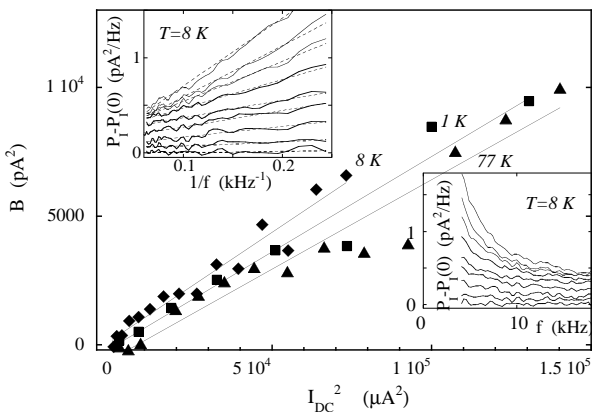


Fig. 4. $1/f$ noise in Rope A. Proportionality constant between the low frequency excess noise and $(1/f)$ versus I_{DC}^2 , at different temperatures. Lower inset : Low frequency excess noise $P_I(I_{DC}) - P_I(O)$ versus frequency at $T = 8 K$. Upper inset : Low frequency excess noise versus $1/f$ at $T = 8 K$. The solid lines are linear fits for $f < 16 kHz$ and correspond to $B = 10^{-13}$.

which is quite unphysical at 1 K and at low current. Observation of proximity induced superconductivity in the same type of samples is a confirmation of long inelastic scattering lengths in the tubes and absence of inelastic scattering at the contacts. Of course, this electron phonon scattering is expected to increase at high currents, as has been invoked to explain the strong non-linearities observed in low temperature transport in individual tubes observed above $I = 5 \mu A$ per tube [23]. We have checked that non linearities for all measured samples discussed here remain extremely small: less than 1 percent of resistance (see inset of fig. 3) and are opposite in sign compared to what is expected for current-induced electron phonon scattering

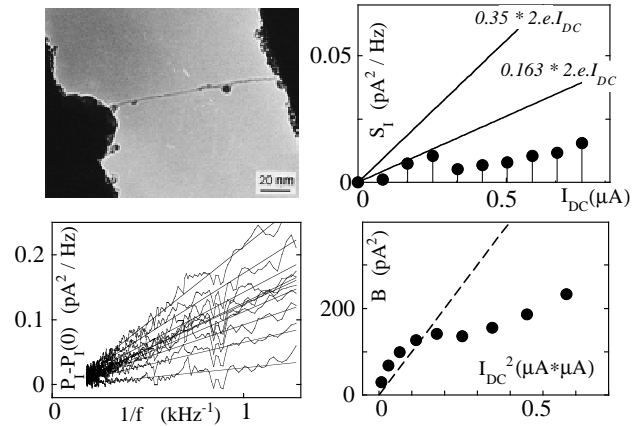


Fig. 5. Individual tube at $T=1 K$. *Upper-left inset:* TEM micrograph of the sample which is $0.15 \mu m$ long. *Upper-right inset:* Shot noise $S_I = P_I(I_{DC}) - P_I(O)$ versus bias current I_{DC} at 1 K for the individual tube. The solid circles correspond to the maximal shot noise compatible with the data, as indicated by the error bars. Note however that the residual background noise of $P_I(O) = 0.04 pA^2/Hz$ measured for this sample is 2.5 times higher than the expected thermodynamic noise at 1 K. The solid lines are the lower and upper bounds for the value of S_I for a two independent channel conductor compatible with the sample resistance $R = 10.5 k\Omega$ measured at room temperature. They correspond to reduction factor of 0.163 and 0.350. *Lower-left inset:* Low frequency excess noise $P_I(I_{DC}) - P_I(O)$ versus $1/f$. The solid lines are linear fits for $f < 5.5 kHz$. *Lower-right inset:* Amplitude of low frequency excess $(1/f)$ noise versus I_{DC}^2 . The solid line is a linear fit only valid in the low current region from which the coefficient $B = 10^{-9}$ is extracted).

where R should exhibit an increase with current instead of the small decrease observed here.

ii) Ballistic transport through the tubes in a Landauer type of picture [25]. Assuming that carbon nanotubes in a rope behave as independent 2 channel conductors characterized by their transmissions t_i^a and t_i^b , the total conductance of a rope reads :

$$G = G_0 \sum_i (t_i^a + t_i^b) \quad (2)$$

where $G_0 = 2e^2/h$ is the conductance quantum. The shot noise spectral density through the sample biased with the voltage V in the limit $eV \gg k_B T$ is:

$$S_I = 2G_0 eV \sum_i [t_i^a (1 - t_i^a) + t_i^b (1 - t_i^b)] \quad (3)$$

From the dimensionless value of the conductance of Rope A, which varies with T between $\mathcal{G} = G/G_0 = 11.80$ and $\mathcal{G} = 11.93$ it is possible to deduce the configuration of t_i which yields the minimum value of S_I . It is obtained assuming that all tubes except 6 are insulator which means $N_G = 12$ conducting channels for the rope. If 5 tubes are perfectly conducting ($t_i^a = t_i^b = 1$) and the sixth one is such that: $t_6^a = 1$ then $t_6^b = \mathcal{G} - \text{int}(\mathcal{G}) = \mathcal{G} - 11$. We deduce a shot noise reduction factor $t_6^b (1 - t_6^b) / \mathcal{G}$ which yields S_I between $eI/70$ and $eI/200$ when the resistance

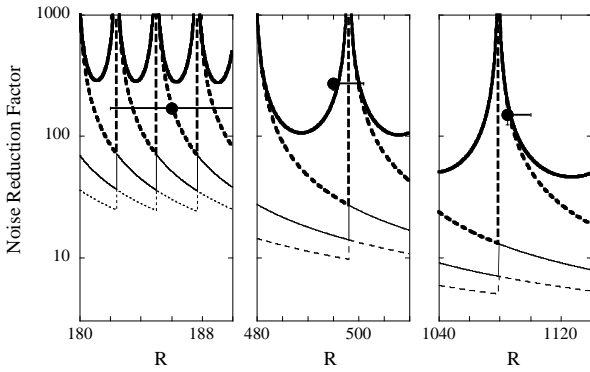


Fig. 6. Noise reduction factor for a multichannel conductor of resistance R (axis) for 2 different types of transmission distribution. Thick solid line: all conducting channels have a full transmission except one. This corresponds to the maximum reduction attainable considering all the combinations of transmissions. Other lines: all conducting channels have the same transmission $t_i = \mathcal{G}/N_G$. Thick dashed line: $N_G = \text{Int}(\mathcal{G}) + 1$ (minimum number of channels compatible with \mathcal{G}), Thin solid line: $N_G = \text{Int}(\mathcal{G}) + 2$, Thin dashed line: $N_G = \text{Int}(\mathcal{G}) + 3$ The full circles represents present data. The horizontal segments represent the temperature dependences of the samples.

varies between 1085Ω and 1100Ω i.e. of the order of the observed value (see fig.6). It is interesting to consider also the situation in which all the conducting channels have the same transmission. The maximum noise reduction is obtained again when the number of conducting channels N_G is minimum ($N_G = \text{Int}(\mathcal{G}) + 1$). Fig.6 (thick dashed line) shows that reduction factor peaks for integer values of \mathcal{G} , which correspond again to a full transmission ($t_i = 1$) for all conducting channels.

It is thus in principle nearly possible to explain the low level of shot noise measured in the ropes by considering them as a small number ($\text{Int}(\mathcal{G}) + 1$ channels) of independent quasi ballistic conductors of non interacting electrons. For Rope A, this implies for all channels either a transmission t_i equal to 0 or 1 within δt_i ($|\delta t_i| < 0.1$), with the very strong additional constraint that $\sum_i |\delta t_i| < 0.1$. The preliminary results obtained on the individual tube presented in fig.5 also indicate that the level of shot noise is reduced by more than a factor 3 compared to the maximum shot noise reduction obtained assuming that the nanotube consists of 2 independent channels of non interacting electrons. As discussed above the presence of electron-electron interactions can induce a Luttinger liquid state in a 1D system resulting in non-integer charge excitations with $e^* = ge$ where g is the LL parameter estimated to be of the order of 0.3 in carbon nanotubes. Replacing e by e^* in (3) could help fit our experimental reduction of shot noise for all the samples. However we want to stress that the assumption we have made, that only a very small number of nearly perfectly ballistic conducting tubes contribute to transport through the ropes, implies that the current per tube in our experiments can exceed $60 \mu A$ per tube within a rope without damaging the sample nor giving rise to observable non-

linearities. This is in apparent contradiction with previous experiments both in individual tubes and ropes where non-linearities start to show up at $5 \mu A$ and the maximum current which can be sustained by an individual SWNT is observed to be of the order of $20 \mu A$ [23, 26]. It is possible that the absence of non-linearities in our samples is due to their suspended character: emitted phonons can not leave the tubes via the substrate and are reabsorbed by other electrons as observed in point contacts [27] which can also sustain very high currents.

Assuming instead that more than $\text{Int}(\mathcal{G}) + 1$ channels participate to transport will yield a shot noise reduction factor which decreases very rapidly with $\text{Int}(\mathcal{G}) + 1$. For example, let us assume that N_G conducting channels have the same transmission \mathcal{G}/N_G . Fig.6 (thin lines) shows the noise reduction factor for $N_G = \text{Int}(\mathcal{G}) + 2$ and $N_G = \text{Int}(\mathcal{G}) + 3$. For all ropes, it falls well below the experimental values. For $N_G = 2 * \text{Int}(\mathcal{G})$, the effect is much more drastic with a reduction factor approaching 2. Such low reductions cannot explain our results. It therefore seems in this hypothesis that a much smaller effective charge is needed to explain the severe reduction we observe.

It may then be necessary to take into account strong correlations between the current noise in the different tubes constituting the rope or coupling between the two channels within each nanotube. The effect of inter-tubes correlations has been recently discussed theoretically by Trauzettel *et al.* [10] in a more sophisticated experimental setup where coupled nanotubes can be independently addressed. These authors propose a scenario with a perfect locking of the noise in different tubes. Such positive correlations are expected to increase shot noise. Our experiments indicate instead strong anticorrelations.

Note that strong Coulomb repulsion between electrons in vacuum diodes or triodes have already been shown 60 years ago [28] to provide a very efficient mechanism for shot noise reduction. At this stage both more experimental and theoretical work are needed to understand our results.

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References

1. M.S. Dresselhaus, G. Dresselhaus, and P.C. Eklund, *Science of Fullerenes and Carbon nanotubes* (Academic, San Diego, 1996).
2. N. Hamada, S. I. Sawada and A. Oshiyama, Phys. Rev. Lett. **68**, 1579 (1992).
3. C. Kane, L. Balents and M. P. Fisher, Phys. Rev. Lett **79**, 5086 (1997).
4. R. Egger, A. Gogolin, Phys. Rev. Lett. **79**, 5082 (1997) ; R. Egger, Phys. Rev. Lett. **83**, 5547 (1999).
5. M. Bockrath *et al.*, Nature **397**, 598 (1999).

6. Ya.M. Blanter and M. Büttiker Physics Reports **3361**, (2000).
7. C.L. Kane and M.P.A. Fisher, Phys. Rev. Lett.**72**,724 (1994). K.-V. Pham, M. Gabay and P. Lederer, Phys. Rev. B **61**, 16397 (2000).
8. R. de Picciotto *et al.*. Nature **389**, 162 (1997); L. Saminadayar, D.C. Glattli, Y. Jin and B. Etienne Phys. Rev. Lett.**79**,2526 (1997).
9. V.V. Ponomarenko and N. Nagaosa, Phys.Rev.B **60**,16865 (1999).
10. B. Trauzettel, R. Egger and H. Grabert cond-mat/0109022.
11. K.I. Imura, K.V. Pham, P. Lederer and F. Piechon cond-mat/0111358 and to be published.
12. C. Journet, *et al.*, Nature **388**, 756 (1997).
13. L. Vaccarini, *et al.*, C.R. Acad. Sci.**327**, 925 (1999).
14. M. Burghard *et al.*, *Electronic Properties of Novel Materials* (Eds H. Kuzmany *et al.*.) AIP, New York (1998).
15. A.Yu. Kasumov, I.I. Khodos, P.M. Ajayan, C. Colliex, Europhys. Lett. **34**, 429 (1996); A.Yu. Kasumov *et al.*, Europhys. Lett. **43**, 89 (1998).
16. typically a Cr/Au or Pt/Au bilayer.
17. A.Yu. Kasumov *et al.*, Science **284**, 1508 (1999).
18. J.E. Fisher *et al.*. Phys. Rev. **B,55**, 4921 (1997).
19. P. Dutta and P.M. Horn, Rev. Mod. Phys. **53**,497 (1981).
20. K. Ralls and R. Burhman, Phys. Rev. Lett. **60**,2434 (1988).
21. P.G. Collins *et al.* Appl.Phys.Lett. **76**, 89 4 (2000),L. Roschier *et al.*. Appl. Phys. Lett. **78**, 3295 (2001).
22. H.W. Postma *et al.*. to appear in Moriond 2001 proceedings.
23. Z. Yao, C.L. Kane, C. Decker Phys. Rev. Lett. **84**, 2941 (2000).
24. B.J. Van Wees *et al.*. Phys. Rev. Lett.**60**,848 (1988).
25. R. Landauer and Th. Martin Physica B **175**,167(1991).
26. P. Collins *et al.*. Science **292**,706 (2001).
27. O.P. Balakahin and I.I. Kulik, Physica **B 218**, 50 (1996).
28. A.W. Hull and N.H. Williams, Phys. Rev. **25**, 147 (1925); A. J. Rack, Bell Syst. Tech. J. **17**, 592 (1938);D. O. North, RCA Rev.**4**, 441 (1940)